

# Multifractal formalisms for oscillating singularities

Stéphane Jaffard

Joint work with Patrice Abry and Stéphane Roux (ENS de Lyon)

Wavelets and Fractals  
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Signals whose instantaneous frequency changes

Examples :

- ▶ Ultrasonic signals emitted by bats
- ▶ Gravitational waves
- ▶ Structures predicted in fully developed turbulence
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Local singular behaviors of signals

(B. Torresani, Y. Meyer)

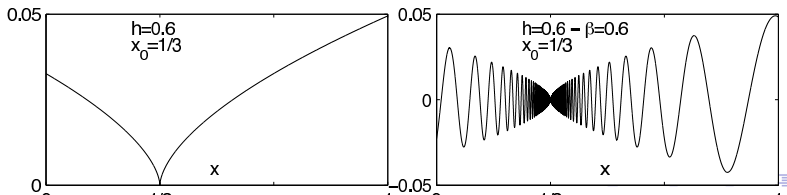
$$f(x) - f(x_0) \sim |x - x_0|^\alpha \sin \left( \frac{1}{|x - x_0|^\beta} \right)$$

# Types of pointwise singularities

**Cusps** :  $f(x) - f(x_0) = |x - x_0|^\alpha$

**Chirps** :  $f(x) - f(x_0) = |x - x_0|^\alpha \sin \left( \frac{1}{|x - x_0|^\beta} \right)$

We look for this typical behavior (not the exact form)



# Spectrum of singularities

The Hölder exponent of  $f$  at  $x_0$  is

$$h_f(x_0) = \sup\{\alpha : f \in C^\alpha(x_0)\}.$$

The Iso-Hölder sets of  $f$  are

$$E_H = \{x_0 : h_f(x_0) = H\}.$$

The spectrum of singularities of a function  $f$  is

$$d_f(H) = \dim(E_H)$$

where  $\dim$  stands for the Hausdorff dimension.

# Wavelet techniques

(initiated by A. Arneodo et al.)

A **wavelet basis** on  $\mathbb{R}$  is generated by a **smooth, well localized, oscillating** function  $\psi$  such that the

$$2^{j/2}\psi(2^jx - k), \quad j, k \in \mathbb{Z}$$

form an orthonormal basis of  $L^2(\mathbb{R})$ .

## Notations :

**Dyadic intervals :**  $\lambda = \left[ \frac{k}{2^j}, \frac{k+1}{2^j} \right[$

**Wavelet coefficients :**  $c_{j,k} = c_\lambda = 2^j \int f(x)\psi(2^jx - k)dx$

If  $\psi_\lambda(x) = \psi(2^jx - k)$ , then  $f(x) = \sum_{\lambda} c_\lambda \psi_\lambda(x)$ .

# Wavelet leaders

Let  $\lambda$  be a dyadic cube ;  $3\lambda$  is the cube of same center and three times wider.

Let  $f$  be a **bounded function** ; the **wavelet leaders** of  $f$  are the quantities

$$L_\lambda = \sup_{\lambda' \subset 3\lambda} |c_{\lambda'}|$$

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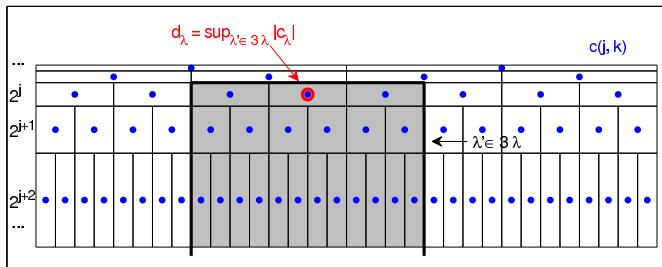
$$L_\lambda = \sup_{\lambda' \subset 3\lambda} |c_{\lambda'}|$$

**Notations :** Let  $x_0 \in \mathbb{R}^d$

$\lambda_j(x_0)$  is the dyadic cube of width  $2^{-j}$  which contains  $x_0$

$$L_j(x_0) = L_{\lambda_j(x_0)} = \sup_{\lambda' \subset 3\lambda_j(x_0)} |c_{\lambda'}|.$$

# Wavelet leaders



A function  $f$  is **uniform Hölder** if  $f \in C^\varepsilon$  for an  $\varepsilon > 0$ , i.e.

$$\exists C > 0 : \quad \forall j, \quad |c_\lambda| \leq C \cdot 2^{-\varepsilon j}.$$

**Characterization of pointwise smoothness** : If  $f$  is uniform Hölder, then

$$\forall x_0 \in \mathbb{R}^d \quad h_f(x_0) = \liminf_{j \rightarrow +\infty} \frac{\log(L_j(x_0))}{\log(2^{-j})}.$$

$$L_j(x_0) \sim 2^{-h_f(x_0)j}$$

# Multifractal formalism

( G. Parisi, U. Frisch, A. Arneodo, S. Jaffard, ....)

$$\Lambda_j = \{\lambda : |\lambda| = 2^{-j}\}$$

$$T_{p,j} = 2^{-dj} \sum_{\lambda \in \Lambda_j} (d_\lambda)^p \sim 2^{-\eta_f(p)j}$$

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The **Legendre spectrum** of  $f$  is

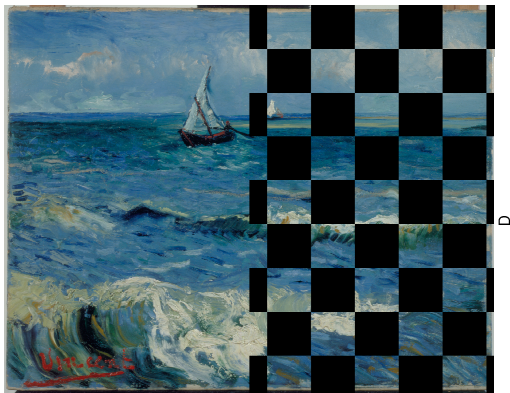
$$\mathcal{L}_f(H) = \inf_{p \in \mathbb{R}} (d + Hp - \eta_f(p))$$

The **wavelet leaders multifractal formalism** holds if

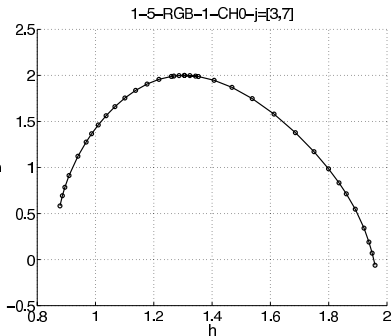
$$d_f(H) = \inf_{p \in \mathbb{R}} (d + Hp - \eta_f(p))$$

# Multifractal analysis of paintings : The Van Gogh challenge

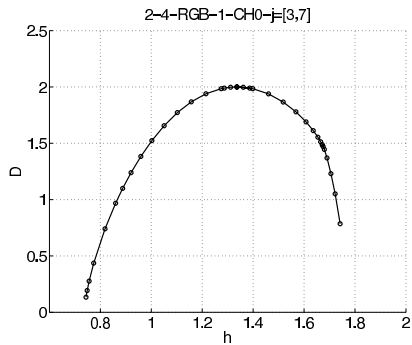
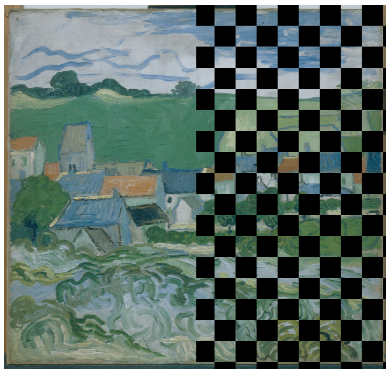
Collaboration with D. Rockmore (Dartmouth) and H. Wendt (Purdue)

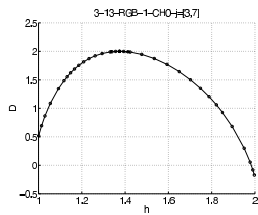
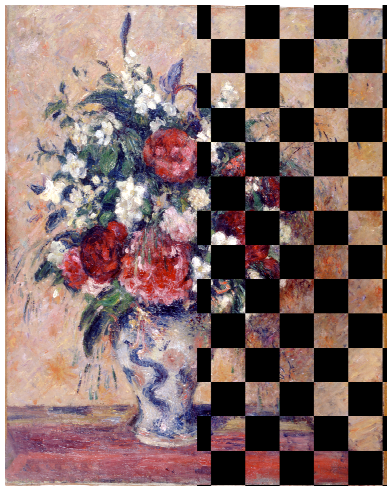


Van Gogh (f415) Arles -Saint Rémy



► f799 (Van Gogh)





Unknown (f248a)

# Princeton Experiment

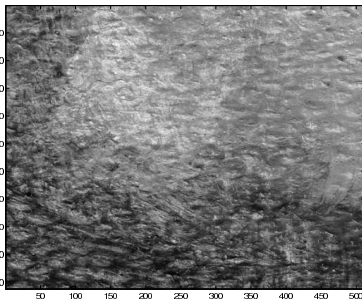
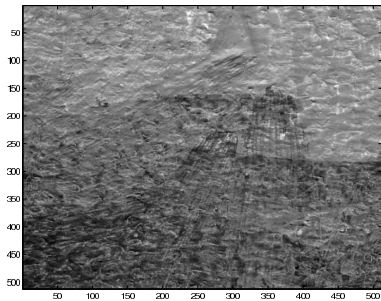
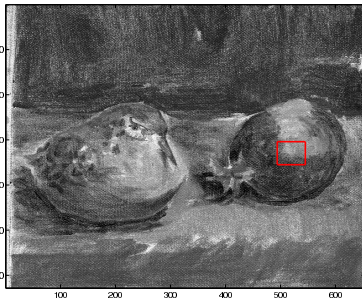
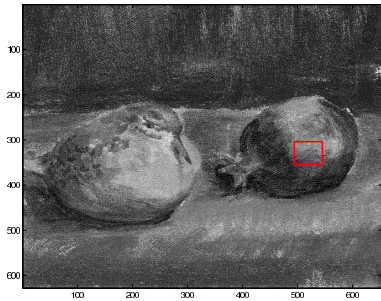
- Experiment design :
  - same Painter (Charlotte Casper) does Original and Copies
  - a series of 7 small paintings,
  - different set of materials (various brushes, grounds, paints)
  - Original and Copies with same set of materials
  - high resolution digitalisation, under uniform conditions.
- Description :

| Pair | Ground            | Paint    | Brushes |
|------|-------------------|----------|---------|
| 1    | CP Canvas         | Oils     | S& H    |
| 2    | CP Canvas         | Acrylics | S& H    |
| 3    | Smooth CP Board   | Oils     | S& H    |
| 4    | Bare linen canvas | Oils     | S       |
| 5    | Chalk and Glue    | Oils     | S       |
| 6    | CP Canvas         | Acrylics | S       |
| 7    | Smooth CP Board   | Oils     | S       |

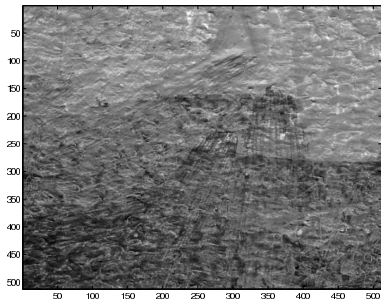
## Charlotte2's Original & Copy



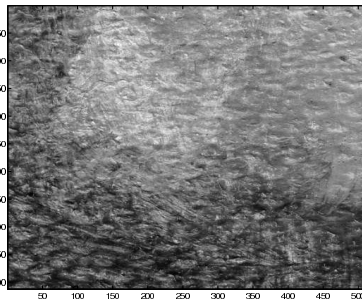
# Charlotte2 MF



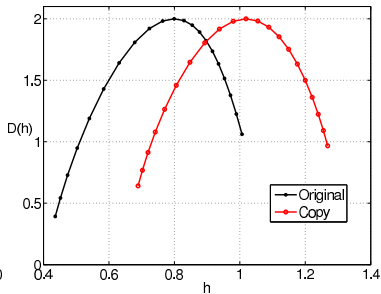
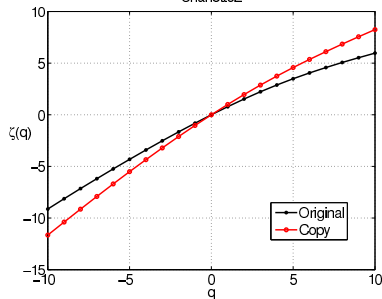
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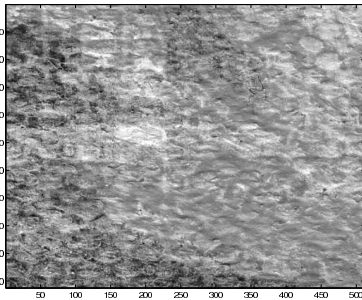
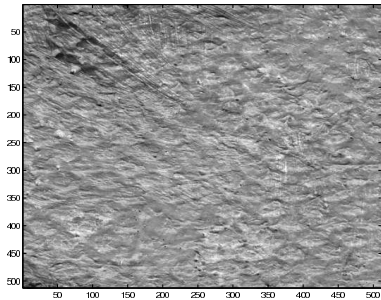
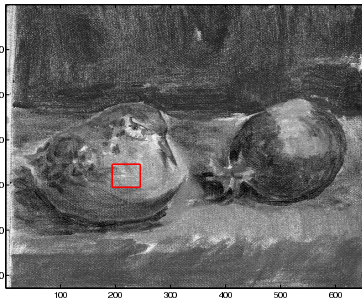
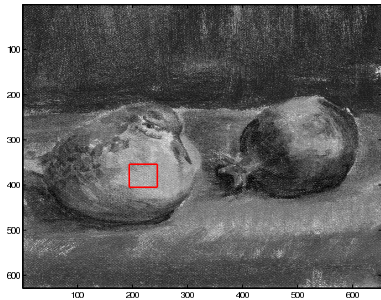
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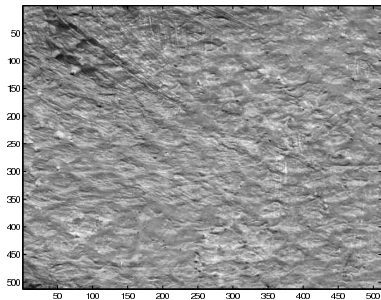
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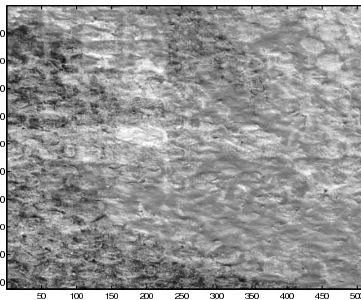
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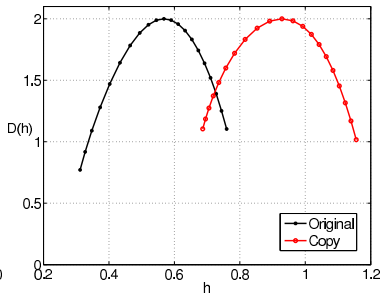
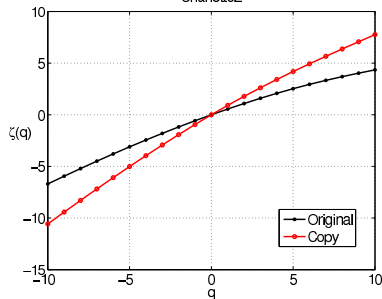
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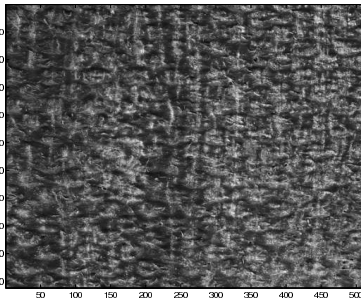
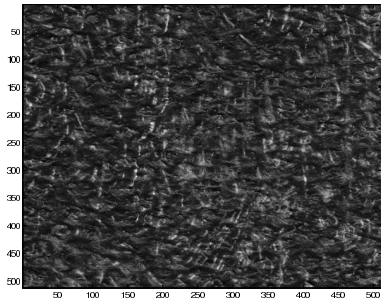
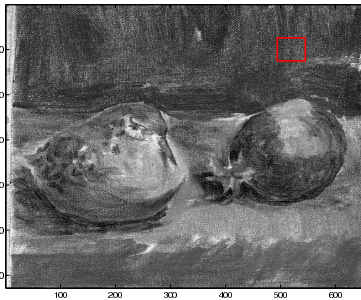
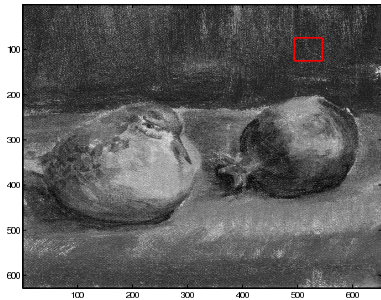
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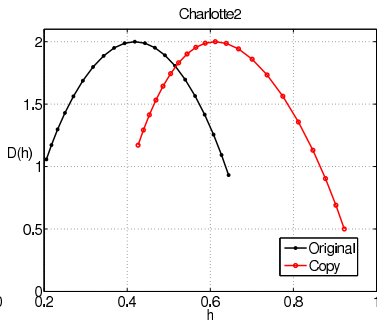
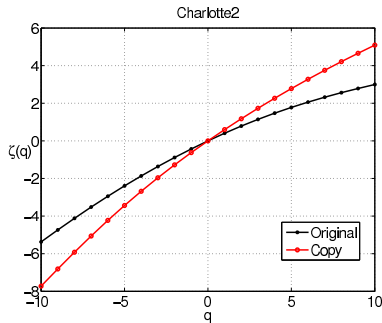
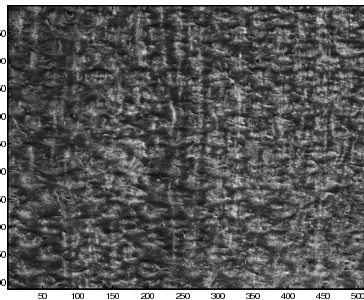
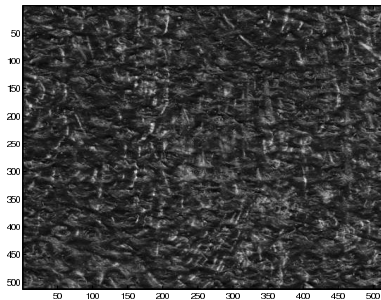
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# Multifractal analysis of oscillating singularities

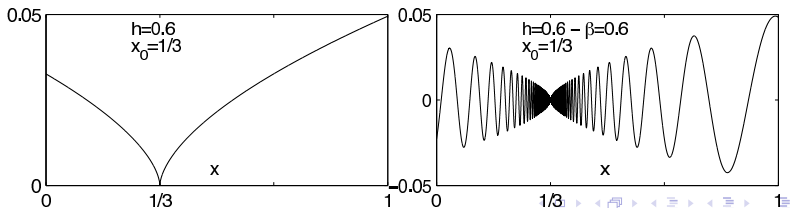
**Cusps** :  $f(x) - f(x_0) = |x - x_0|^H$

After one integration :  $f^{(-1)}(x) - f^{(-1)}(x_0) \sim \frac{1}{H} |x - x_0|^{H+1}$

**Chirps** :  $f(x) - f(x_0) = |x - x_0|^H \sin \left( \frac{1}{|x - x_0|^\beta} \right)$

After one integration :

$$f^{(-1)}(x) - f^{(-1)}(x_0) = \frac{|x - x_0|^{H+(1+\beta)}}{\beta} \cos \left( \frac{1}{|x - x_0|^\beta} \right) + \dots$$



# Fractional Integration

The **fractional integral of order  $s$**  is the operator  $\mathcal{I}^s$  satisfying

$$\widehat{\mathcal{I}^s f}(\xi) = (1 + |\xi|^2)^{-s/2} \hat{f}(\xi)$$

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- ▶ If  $f$  is a cusp at  $x_0$ , then  $h_{\mathcal{I}^s f}(x_0) = h_f(x_0) + s$
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**Definition** : Let  $f : \mathbb{R} \rightarrow \mathbb{R}$  be a locally bounded function.  
The **oscillation exponent** of  $f$  at  $x_0$  is

$$\beta_f(x_0) = \left( \frac{\partial(h_{\mathcal{I}^s f}(x_0))}{\partial s} \right)_{s=0^+} - 1$$

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$f$  has a **cusp** at  $x_0$  if  $\beta_f(x_0) = 0$ . It follows that

$$h_{\mathcal{I}^s f}(x_0) = h_f(x_0) + s$$

# Pseudo-fractional integration

Heuristic :

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**Theorem :** If  $f$  is uniform Hölder, then  $\mathcal{I}^s(f)$  and  $\mathcal{J}^s(f)$  share the same pointwise exponents and spectra

Algorithm :

- ▶ Operate pseudo-fractional integration of order  $s$
- ▶ Perform the multifractal analysis of this new function

One obtains the **integrated Legendre spectrum** :  $\mathcal{L}_f^s(H) = \mathcal{L}_{\mathcal{I}^s(f)}^s(H)$

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If all points are cusps, then the integrated spectrums of singularities are shifted

$$d_f^s(H) = d_f(H - s)$$

**Heuristics :**

If all points are cusps, then the integrated Legendre spectrums are shifted

# Fractionally integrated spectra

**Definition** : A uniformly Hölder function  $f$  is of cusp type

$$\forall \lambda, \exists \lambda' \subset \lambda \quad j' = j + o(j) \text{ et } |c_{\lambda'}| \geq L_{\lambda} 2^{-o(j)}$$

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**Proposition** : Si  $f$  is uniformly of cusp type then the integrated Legendre spectra satisfy

$$\mathcal{L}_f^s(H) := \mathcal{L}_{\mathcal{I}^s f}(h) = \mathcal{L}_f(H - s)$$

**Exemples** : FBM, Random Wavelet cascades of A. Arneodo, The measure-based random wavelet series of J. Barral and S. Seuret,...

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**Exemples** : FBM, Random Wavelet cascades of A. Arneodo, The measure-based random wavelet series of J. Barral and S. Seuret,...

**Heuristics** :

Cusps= Signature of clustering of large wavelet coefficients

Oscillating singularities= Dispersion of large wavelet coefficients

**Problem** : Oscillating singularities can be present even if the integrated Legendre spectra are shifted

# The $\beta$ -spectrum

The **iso-oscillating sets** are

$$F_\beta = \{x : \beta_f(x) = \beta\}$$

The  **$\beta$ -spectrum** is

$$\mathcal{D}_f(\beta) = \dim(F_\beta)$$

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**Heuristic :** Let  $s$  “small enough ” be given. If  $f$  has an oscillating singularity of exponents  $(H, \beta)$  at  $x_0$ , then

$$L_{\lambda_j(x_0)}(f)(x_0) \sim 2^{-Hj} \quad \text{and} \quad L_{\lambda_j(x_0)}(\mathcal{I}^s f) \sim 2^{-(H+s(1+\beta))j}$$

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The  **$\beta$ -leaders** by

$$B_\lambda = 2^j \left( \frac{L_{\lambda_j(x_0)}(\mathcal{I}^s f(x_0))}{L_{\lambda_j(x_0)}(f)(x_0)} \right)^{1/s} \sim 2^{-\beta j}$$

# The $\beta$ -formalism

The  $\beta$ -structure function is  $B_{p,j} = 2^{-dj} \sum_{\lambda \in \Lambda_j} (B_\lambda)^p$

The  $\beta$ -scaling function is

$$\omega_f(p) = \liminf_{j \rightarrow +\infty} \frac{\log(B_{p,j})}{\log(2^{-j})}$$

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The  $\beta$ -formalism holds if

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$$\mathcal{D}_f(\beta) = \inf_{p \in \mathbb{R}} (d + \beta p - \omega_f(p))$$

**Heuristic :** If  $\mathcal{D}_f(\beta)$  is supported by a point, then one expects that there are no oscillating singularities

# The $\beta$ -formalism

The  $\beta$ -structure function is  $B_{p,j} = 2^{-dj} \sum_{\lambda \in \Lambda_j} (B_\lambda)^p$

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**Theorem :**

- ▶ If  $f$  is of cusp-type, then  $\mathcal{D}_f(\beta)$  is supported by a point
- ▶ The  $\beta$ -formalism holds for lacunary wavelet series

# The grandcanonical formalism

## Spectrum of oscillating singularities

$$\mathbb{D}_f(H, \beta) = \dim(\{x_0 : h_f(x_0) = H \text{ et } \beta_f(x_0) = \beta\})$$

# The grandcanonical formalism

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If  $f$  has an oscillating singularity of exponents  $(H, \beta)$  at  $x_0$ , then

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The grandcanonical structure function is

$$G_{p,q,j} = 2^{-dj} \sum_{\lambda \in \Lambda_j} (L_\lambda)^p (B_\lambda)^q$$

The grandcanonical scaling function is

$$\forall p, q \in \mathbb{R}, \quad \Omega_f(p, q) = \liminf_{s \rightarrow 0} \liminf_{j \rightarrow +\infty} \frac{\log(G_{p,q,j})}{\log(2^{-j})}$$

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The grandcanonical multifractal formalism holds if

$$\mathbb{D}_f(H, \beta) = \inf_{p, q \in \mathbb{R} \times \mathbb{R}} (d + Hp + \beta q - \Omega_f(p, q))$$

## Theorem :

The grandcanonical multifractal formalism holds for cusp-type functions and for lacunary wavelet series

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## Challenges :

## Turbulence

Oscillating singularities on sets of small dimension

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## Challenges :

### Turbulence

Oscillating singularities on sets of small dimension

### Alternative stochastic processes which display oscillating singularities

**Heuristic :** oscillating singularities are present when wavelet coefficients are both **sparse** and **dispersed** without interscale correlations